



Condensed Matter Physics

Crystal structure

Istvan Groma

ELTE

October 1, 2018





Periodic systems



Primitive lattice vectors

$$\underline{a}_1, \quad \underline{a}_2, \quad \underline{a}_3$$



Periodic systems

Primitive lattice vectors

$$\underline{a}_1, \quad \underline{a}_2, \quad \underline{a}_3$$



Bravais lattice

$$\underline{R}_n = n_1 \underline{a}_1 + n_2 \underline{a}_2 + n_3 \underline{a}_3$$

Periodic systems

Primitive lattice vectors

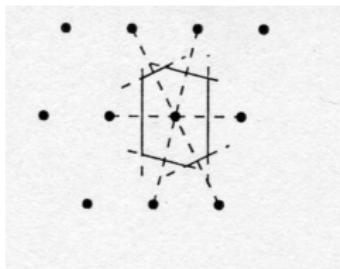
$$\underline{a}_1, \quad \underline{a}_2, \quad \underline{a}_3$$



Bravais lattice

$$\underline{R}_n = n_1 \underline{a}_1 + n_2 \underline{a}_2 + n_3 \underline{a}_3$$

Wigner-Seitz cell





Points and directions

Bravais lattice points

$$\underline{r} = x_1 \underline{a}_1 + x_2 \underline{a}_2 + x_3 \underline{a}_3$$

$$\left(\frac{1}{2}, \frac{1}{2}, 0 \right)$$



Points and directions

Bravais lattice points

$$\underline{r} = x_1 \underline{a}_1 + x_2 \underline{a}_2 + x_3 \underline{a}_3$$

$$\left(\frac{1}{2}, \frac{1}{2}, 0 \right)$$

Directions

$$u \underline{a}_1 + v \underline{a}_2 + w \underline{a}_3$$

$$[u, v, w]$$



Points and directions

Bravais lattice points

$$\underline{r} = x_1 \underline{a}_1 + x_2 \underline{a}_2 + x_3 \underline{a}_3$$

$$\left(\frac{1}{2}, \frac{1}{2}, 0 \right)$$

Directions

$$u \underline{a}_1 + v \underline{a}_2 + w \underline{a}_3$$

$$[u, v, w]$$

Reciprocal lattice

$$\underline{a}_{1,2,3} \rightarrow \underline{b}_{1,2,3} \quad \underline{a}_i \cdot \underline{b}_j = 2\pi \delta_{ij}$$

$$\underline{b}_1 = 2\pi \frac{\underline{a}_2 \times \underline{a}_3}{\det} \quad \underline{b}_2 = 2\pi \frac{\underline{a}_3 \times \underline{a}_1}{\det} \quad \underline{b}_3 = 2\pi \frac{\underline{a}_1 \times \underline{a}_2}{\det}$$



Points and directions

$$\begin{pmatrix} a_{1x} & a_{2x} & a_{3x} \\ a_{1y} & a_{2y} & a_{3y} \\ a_{1z} & a_{2z} & a_{3z} \end{pmatrix} \begin{pmatrix} b_{1x} & b_{1y} & b_{1z} \\ b_{2x} & b_{2y} & b_{2z} \\ b_{3x} & b_{3y} & b_{3z} \end{pmatrix} = \begin{pmatrix} 2\pi & & \\ & 2\pi & \\ & & 2\pi \end{pmatrix}$$

$$v_r = \frac{(2\pi)^3}{v_c}$$

Brillouin zone = reciprocal Wigner Seitz cell



Points and directions

$$\begin{pmatrix} a_{1x} & a_{2x} & a_{3x} \\ a_{1y} & a_{2y} & a_{3y} \\ a_{1z} & a_{2z} & a_{3z} \end{pmatrix} \begin{pmatrix} b_{1x} & b_{1y} & b_{1z} \\ b_{2x} & b_{2y} & b_{2z} \\ b_{3x} & b_{3y} & b_{3z} \end{pmatrix} = \begin{pmatrix} 2\pi & & \\ & 2\pi & \\ & & 2\pi \end{pmatrix}$$

$$v_r = \frac{(2\pi)^3}{v_c}$$

Brillouin zone = reciprocal Wigner Seitz cell

Miller index

$$\underline{G}_{hkl} = h\underline{b}_1 + k\underline{b}_2 + l\underline{b}_3 \quad (h, k, l)$$



Points and directions

$$\begin{pmatrix} a_{1x} & a_{2x} & a_{3x} \\ a_{1y} & a_{2y} & a_{3y} \\ a_{1z} & a_{2z} & a_{3z} \end{pmatrix} \begin{pmatrix} b_{1x} & b_{1y} & b_{1z} \\ b_{2x} & b_{2y} & b_{2z} \\ b_{3x} & b_{3y} & b_{3z} \end{pmatrix} = \begin{pmatrix} 2\pi & & \\ & 2\pi & \\ & & 2\pi \end{pmatrix}$$

$$v_r = \frac{(2\pi)^3}{v_c}$$

Brillouin zone = reciprocal Wigner Seitz cell

Miller index

$$\underline{G}_{hkl} = h\underline{b}_1 + k\underline{b}_2 + l\underline{b}_3 \quad (h, k, l)$$

Equation for a plane

$$\underline{r} \underline{G}_{hkl} = \text{const.}$$

Crosses

$$\underline{R}_n = n_1 \underline{a}_1 + n_2 \underline{a}_2 + n_3 \underline{a}_3$$

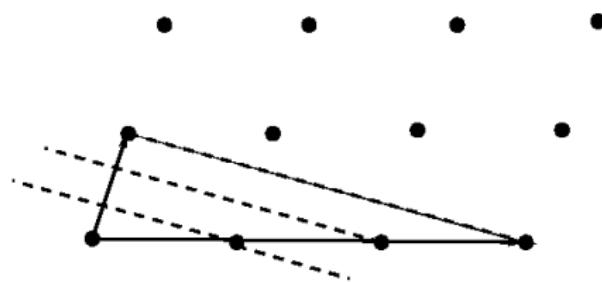
$$\text{const.} = 2\pi m$$

Points and directions

$$h * x_1 + k * x_2 + l * x_3 = m$$

Intercepts

$$\frac{m}{h}, \quad \frac{m}{k}, \quad \frac{m}{l},$$



Distance of the planes

$$d = a_1 \frac{(m+1) - m}{h} \frac{G_{hkl}}{|G_{hkl}|} = \frac{2\pi}{|G_{hkl}|}$$



Rotation of the crystal

Köbös kristály

$$d_{hkl} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$



Rotation of the crystal

Köbös kristály

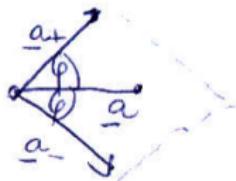
$$d_{hkl} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

Rotation of the crystal

Köbös kristály

$$d_{hkl} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

Rotation



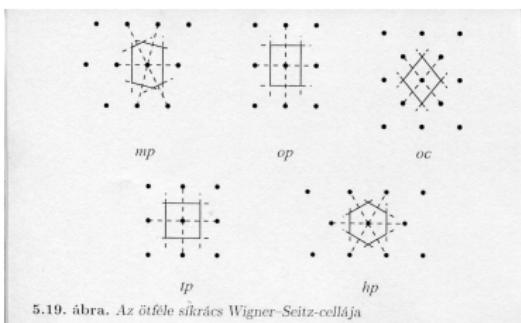
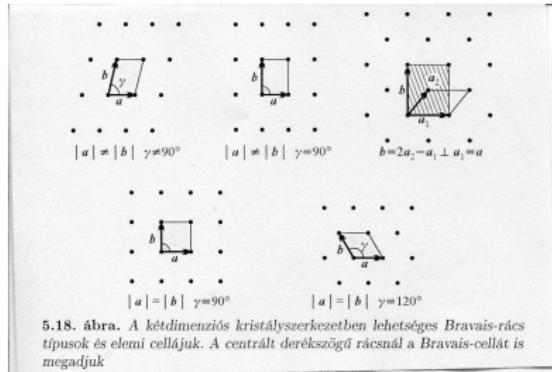
$$\underline{a}_+ + \underline{a}_- = m \underline{a}$$

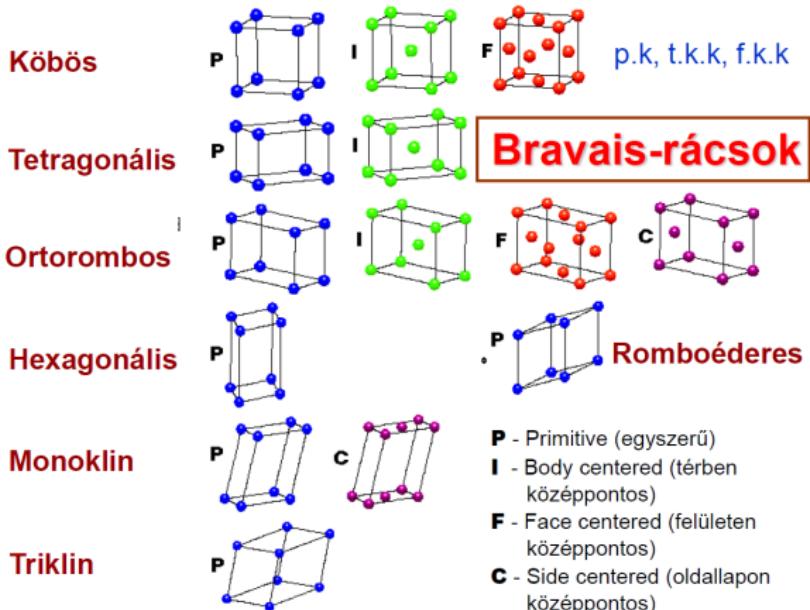
$$\underline{a}_+ + \underline{a}_- = (2 \cos(\varphi), 0, 0) a$$

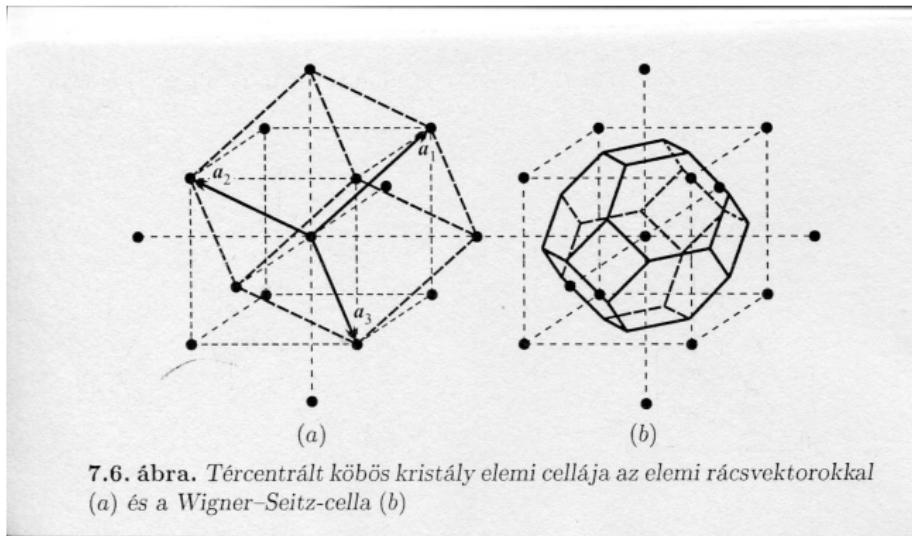
$$\cos(\varphi) = 0, \quad \pm \frac{1}{2}, \quad \pm 1$$

$$\varphi = 0, \quad \pm \frac{2\pi}{2}, \quad \pm \frac{2\pi}{3}, \quad \pm \frac{2\pi}{4}, \quad \pm \frac{2\pi}{6}$$

2D crystals

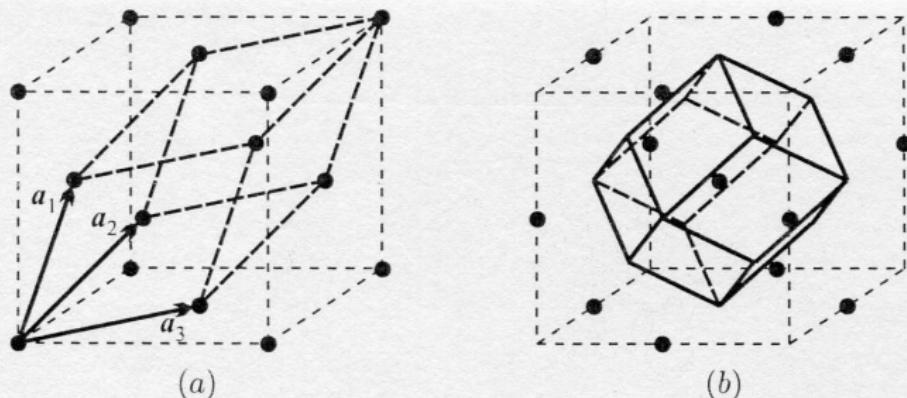






7.6. ábra. Tércentrált köbös kristály elemi cellája az elemi rácsvektorokkal (a) és a Wigner–Seitz-cellá (b)

Cr, Fe, Mo, Nb

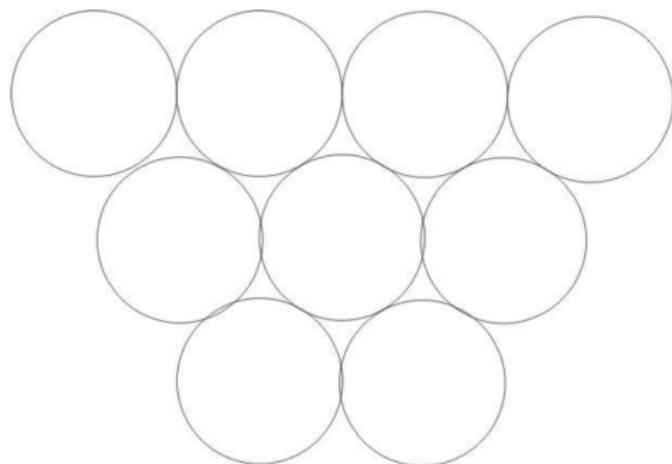


7.9. ábra. Lapcentrált köbös kristály Bravais-cellája az elemi rácsvektorokkal (a) és Wigner–Seitz-cellája (b)

Cu, Ag, Au, Ni

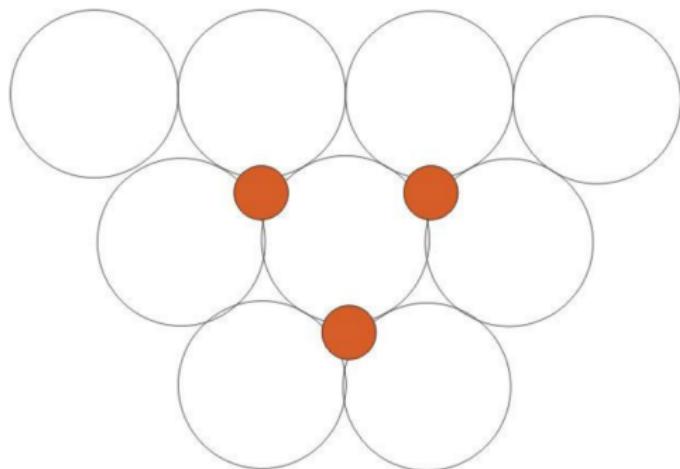


Closed packed crystals



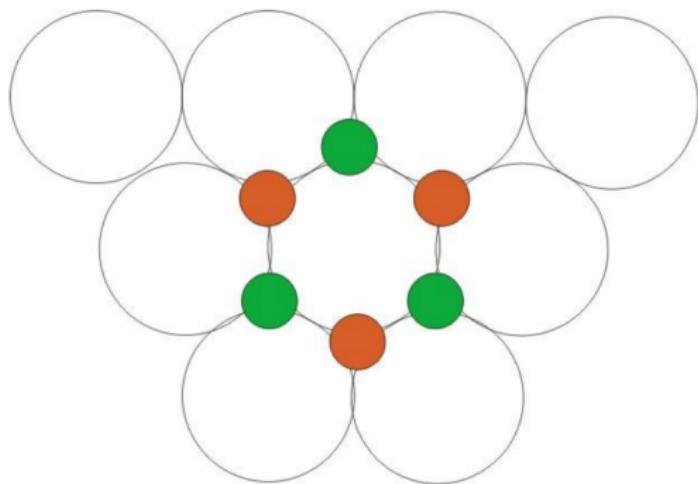


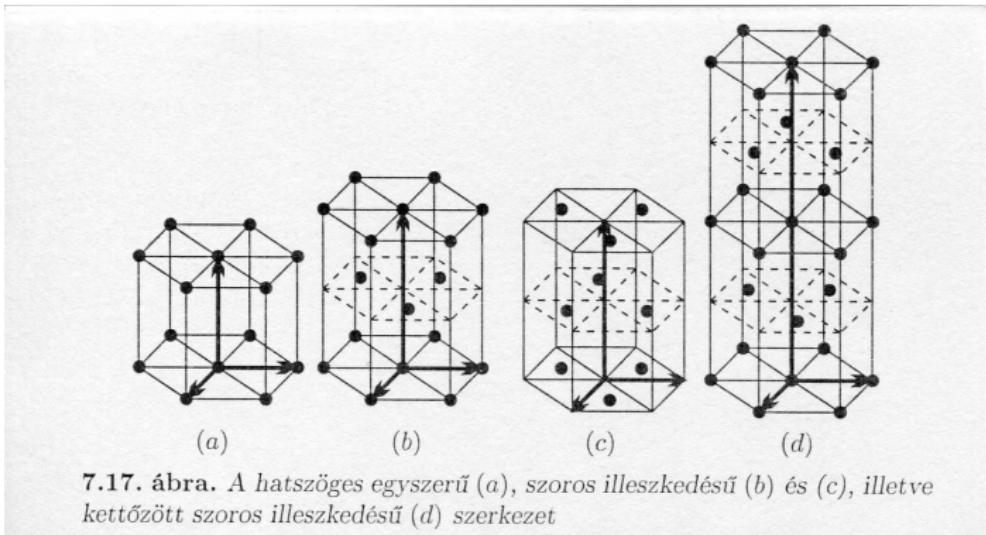
Closed packed crystals





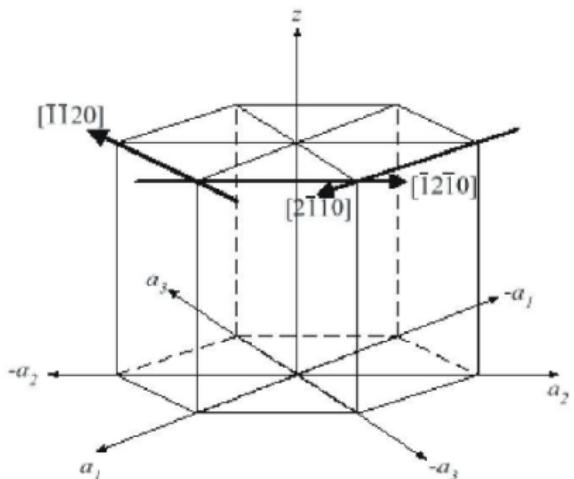
Closed packed crystals



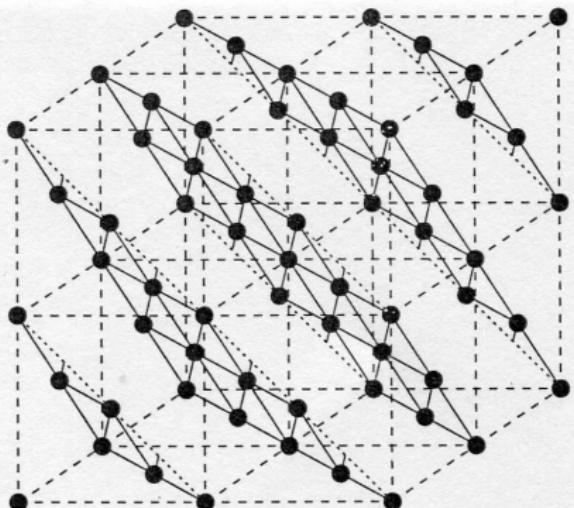


7.17. ábra. A hatszöges egyszerű (a), szoros illeszkedésű (b) és (c), illetve kettőzött szoros illeszkedésű (d) szerkezet

Co, Ti, Zn, Ni, Mg



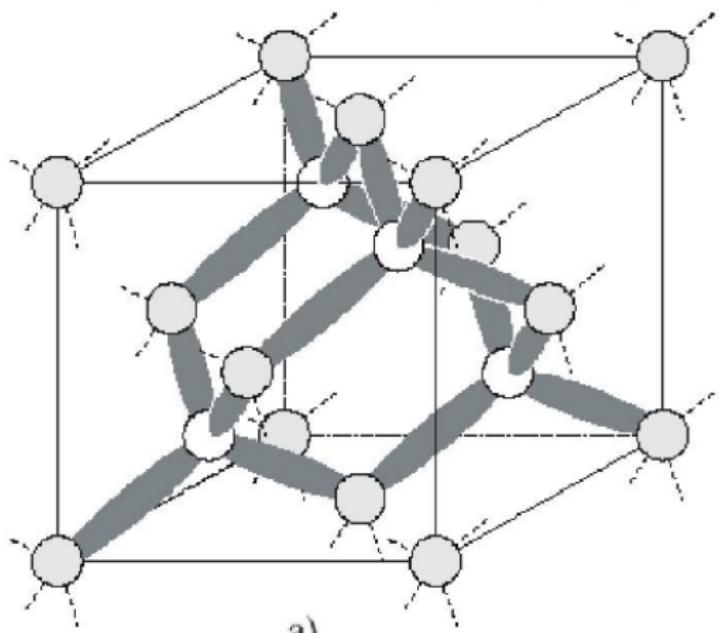
FCC closed packing



7.11. ábra. Az atomok szoros illeszkedésű elhelyezkedése lapcentrált köbös kristályban a térfelületről merőleges síkokban

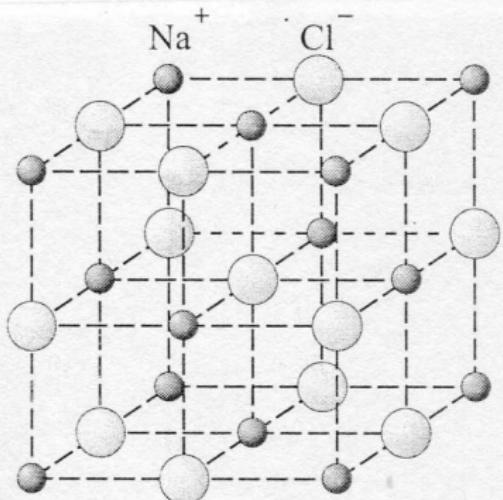


Diamond crystal

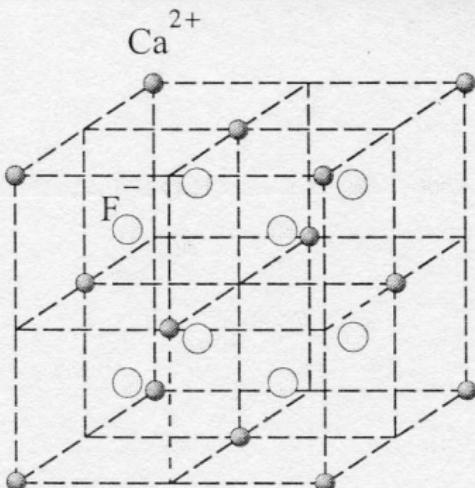




Systems with more atom



(a)



(b)

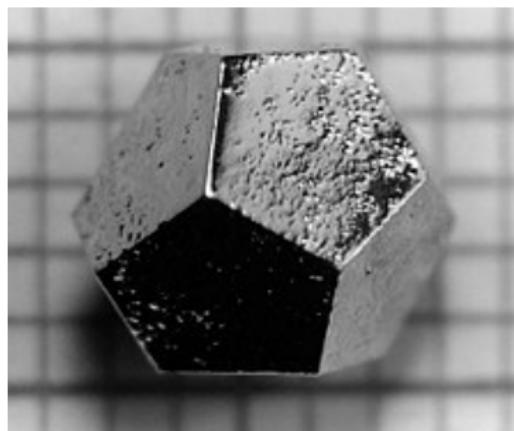
7.13. ábra. Kőszó típusú (B1) (a) és fluorit típusú (C1) (b) kristályszerkezet



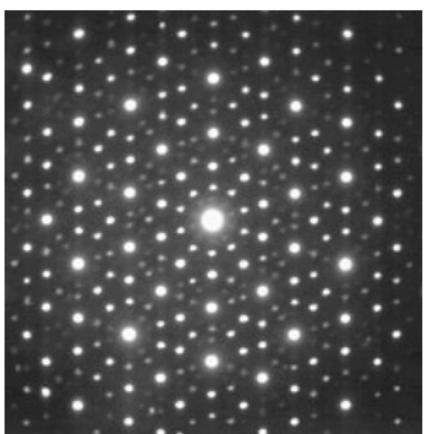
Quasicrystal



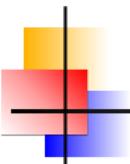
Daniel Dan Shechtman (1982), Nobel price 2011



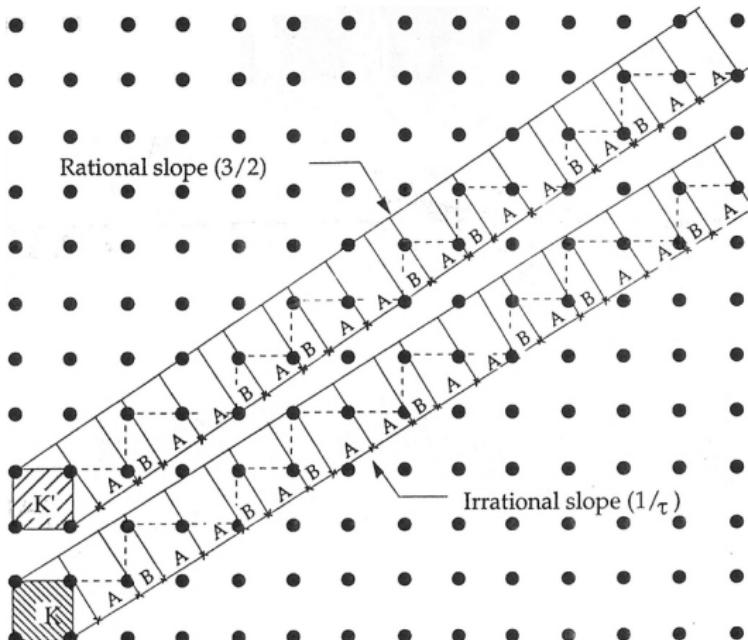
Quasicrystal

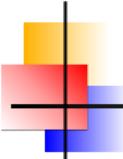


Diffraction image



Construction





Construction

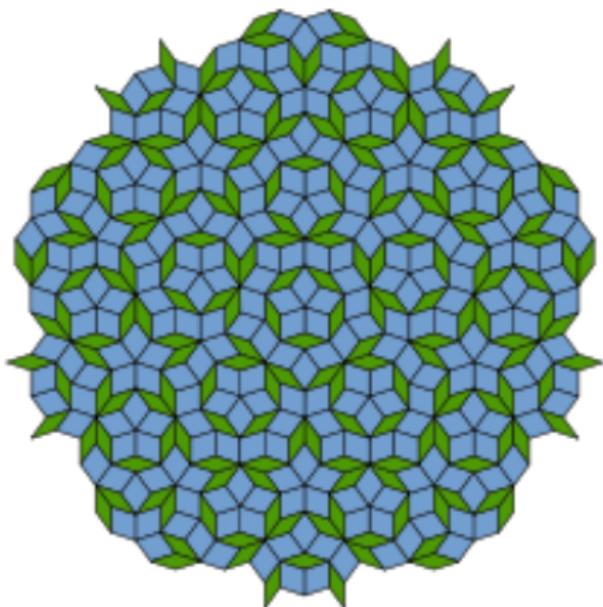


Table 4. The Fibonacci chain.

Fibonacci number	Fibonacci sequence	#L	#S
1	S	0	1
1	L	1	0
2	LS	1	1
3	LSL	2	1
5	LSLLS	3	2
8	LSLLSLSL	5	3
13	LSLLSLSLLSLLS	8	5
21	LSLLSLSLLSLLSLLSLSL	13	8
...



2D Quasicrystal



Penrose tiling

Fat rhombus 72, 72, 108, 108 angles

Slim rhombus 36, 36, 144, 144 angles