



Condensed Matter Physics

Statistical physics of dislocations

István Groma

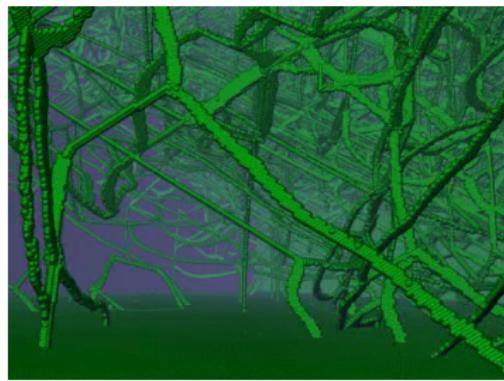
ELTE

May 3, 2019



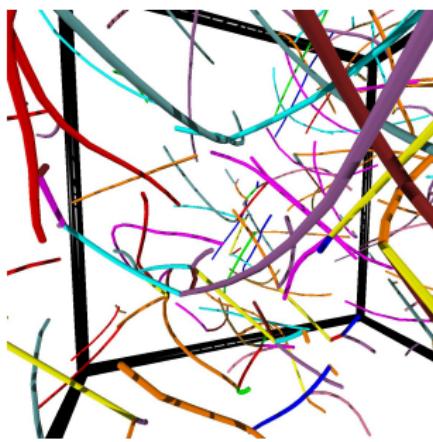
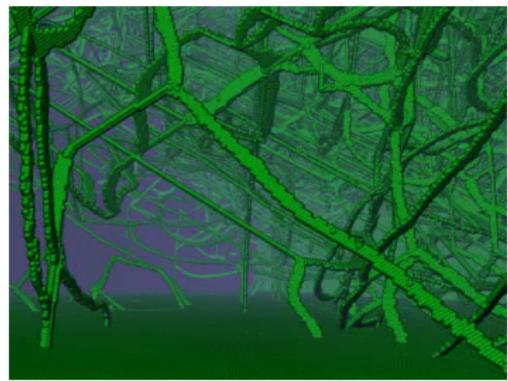


Microscopic picture





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General properties

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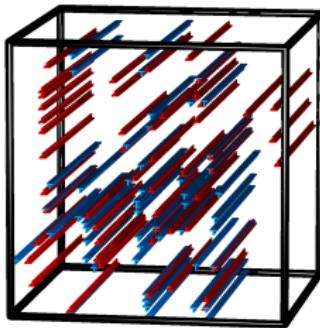


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DDD single slip in 2D

Equation of motion

$$\frac{dx_i}{dt} = M_0 b \left(\sum_{j=1}^N s_i s_j \tau_{ind} (\vec{r}_j - \vec{r}_i) + \tau_{ext} \right)$$

DDD single slip in 2D

Equation of motion with thermal noise

$$\frac{dx_i}{dt} = M' \frac{b}{\sqrt{\rho}} \left(\sum_{j=1}^N s_i s_j \tau_{ind} (\vec{r}_j - \vec{r}_i) + \tau_{ext} \right) + \sqrt{2M' k_B T} \zeta_i(t)$$
$$\langle \zeta_i(t) \zeta_j(0) \rangle = \delta_{ij} \delta(t)$$

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Since

$$\tau_{ind} = \partial_x \partial_y \chi_{ind} \Rightarrow \tau_{ind} = -\partial_x \Phi_{ind}$$

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Since

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$$\frac{dx_i}{dt} = -M' \partial_{x_i} V(\{\vec{r}_i\}) + \sqrt{2M' k_B T} \zeta_i(t)$$

$$V(\{\vec{r}_i\}) = \frac{b}{2\sqrt{\rho}} \sum_{ij} s_i s_j \Phi_{ind} (\vec{r}_i - \vec{r}_j)$$



Fokker–Planck equation

$$\frac{dp(\{x_i\}, t)}{dt} = M' \sum_i \partial_{x_i} [p(\{x_i\}, t) \partial_{x_i} V(\{\vec{r}_i\}) + k_B T \partial_{x_i} p(\{x_i\}, t)]$$

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Steady state solution

$$[\partial_{x_i} V(\{\vec{r}_i\}) + k_B T \partial_{x_i}] p_\infty(\{x_i\}) = 0$$

$$p_\infty(\{x_i\}) = \frac{1}{Z} e^{-\frac{V}{k_B T}}$$

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Time scales

$$t_{relax} = \frac{1}{M' \rho} \frac{\sqrt{\rho}}{b^2 \mu} = \frac{1}{M_0 \rho b^2 \mu} \ll t_T = \frac{1}{M' \rho} \frac{1}{k_B T}$$



Time evolution of the energy

$$\frac{dV}{dt} = \sum_i \vec{v}_i \frac{dV}{d\vec{r}_i} = - \sum_i M' v_i^2 < 0$$



Coarse graining

Set a parallel edge dislocations with $\pm \vec{b}$

$$\nabla^4 \chi = \frac{2\mu}{1-\nu} (b \partial_y) \kappa_d \quad \text{with} \quad \kappa_d = \sum_i s_i \delta(\vec{r} - \vec{r}_i)$$

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Convolution with $w(\vec{r})$

$$\int w(\vec{r} - \vec{r}') \nabla'^4 \chi(\vec{r}') d\vec{r}' = \frac{2\mu}{1-\nu} \int w(\vec{r} - \vec{r}') b \partial_{y'} \kappa_d(\vec{r}) d\vec{r}'$$

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Partial integrations

$$\nabla^4 \int w(\vec{r} - \vec{r}') \chi(\vec{r}') d\vec{r}' = \frac{2\mu}{1-\nu} b \partial_y \int w(\vec{r} - \vec{r}') \kappa_d(\vec{r}') d\vec{r}'$$

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Coarse grained fields

$$\langle \chi \rangle = \int w(\vec{r} - \vec{r}') \chi(\vec{r}') d\vec{r}' \quad \quad \langle \kappa \rangle = \int w(\vec{r} - \vec{r}') \kappa_d(\vec{r}') d\vec{r}'$$

$$\langle \sigma \rangle_{11} = -\partial_{yy} \langle \chi \rangle, \quad \langle \sigma \rangle_{22} = -\partial_{xx} \langle \chi \rangle, \quad \langle \sigma \rangle_{12} = \partial_{xy} \langle \chi \rangle$$

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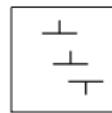
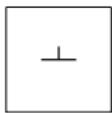
Partial integrations

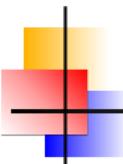
$$\nabla^4 \int w(\vec{r} - \vec{r}') \chi(\vec{r}') d\vec{r}' = \frac{2\mu}{1-\nu} b \partial_y \int w(\vec{r} - \vec{r}') \kappa_d(\vec{r}') d\vec{r}'$$

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Micro - > meso



Equation of motion

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Micro - > meso

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$$-\frac{d}{dt} \rho_{disc}(\vec{r}) = M_0 \vec{b} \frac{d}{d\vec{r}} \left\{ \int \tau_{\text{ind}}(\vec{r}') \rho_{disc}(\vec{r}') \rho_{disc}(\vec{r} - \vec{r}') d\vec{r}' + \tau_{\text{ext}} \rho_{disc}(\vec{r}) \right\},$$

$$\rho_{disc}(\vec{r}, \vec{r}_1, \dots, \vec{r}_N) = \sum_{i=1}^N \delta(\vec{r} - \vec{r}_i)$$

Micro -> meso

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Coarse graining

$$\rho_1(\vec{r}) = <\rho_{disc}(\vec{r})>_{CG}, \quad \rho_2(\vec{r}, \vec{r}') = <\rho_{disc}(\vec{r}') \rho_{disc}(\vec{r} - \vec{r}')>_{CG}$$



Micro - > meso



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$$\frac{\partial \rho_1(\vec{r}_1, t)}{\partial t} + M_0 \vec{b} \frac{d}{d\vec{r}_1} \left[\rho_1(\vec{r}_1, t) \tau_{\text{ext}} + \int \rho_2(\vec{r}_1, \vec{r}_2, t) \tau_{\text{ind}}(\vec{r}_1 - \vec{r}_2) d\vec{r}_2 \right] = 0$$

Micro -> meso

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$$\rho_{disc}(\vec{r}, \vec{r}_1, \dots, \vec{r}_N) = \sum_{i=1}^N \delta(\vec{r} - \vec{r}_i)$$

$$\frac{\partial \rho_+(\vec{r}_1, t)}{\partial t} + M_0 \vec{b} \frac{d}{d\vec{r}_1} \left[\rho_+(\vec{r}_1, t) \tau_{\text{ext}} + \int \{ \rho_{++}(\vec{r}_1, \vec{r}_2, t) - \rho_{+-}(\vec{r}_1, \vec{r}_2, t) \} \tau_{\text{ind}}(\vec{r}_1 - \vec{r}_2) d\vec{r}_2 \right] = 0$$

$$\frac{\partial \rho_-(\vec{r}_1, t)}{\partial t} - M_0 \vec{b} \frac{d}{d\vec{r}_1} \left[\rho_-(\vec{r}_1, t) \tau_{\text{ext}} - \int \{ \rho_{--}(\vec{r}_1, \vec{r}_2, t) - \rho_{-+}(\vec{r}_1, \vec{r}_2, t) \} \tau_{\text{ind}}(\vec{r}_1 - \vec{r}_2) d\vec{r}_2 \right] = 0$$



Plastic shear

Dislocation density tensor

$$\alpha_{31} = b(\rho_+ - \rho_-) = b\kappa \quad \rho = \rho_+ + \rho_-$$



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with $\gamma = \beta_{12}^P$

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Geometrically Necessary Dislocation (GND) density

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Geometrically Necessary Dislocation (GND) density

Deformation rate

$$\frac{\partial \kappa}{\partial t} = -\frac{\vec{b}}{b^2} \frac{d\dot{\gamma}}{dr}.$$



Self-consistent field

Simplest assumption

$$\rho_{ss'}(\vec{r}_1, \vec{r}_2, t) = \rho_s(\vec{r}_1)\rho_{s'}(\vec{r}_2), \quad s, s' \in \{+, -\}.$$

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$$\frac{\partial \rho(\vec{r}, t)}{\partial t} + M_0 \vec{b} \frac{\partial}{\partial \vec{r}} [\kappa(\vec{r}, t) \{\tau_{sc}(\vec{r}, t) + \tau_{ext}\}] = 0$$

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where

$$\tau_{sc}(\vec{r}) = \int \kappa(\vec{r}_1, t) \tau_{ind}(\vec{r} - \vec{r}_1) d\vec{r}_1$$

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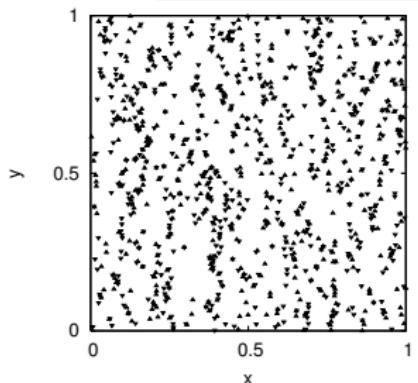
$$\tau_{sc}(\vec{r}) = \int \kappa(\vec{r}_1, t) \tau_{ind}(\vec{r} - \vec{r}_1) d\vec{r}_1$$

τ_{sc} is not "new"

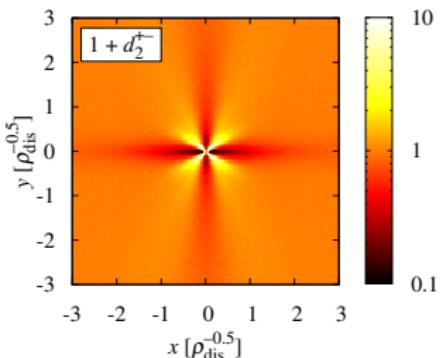
$$\Delta^2 \chi = \frac{2b\mu}{(1-\nu)} \frac{\partial}{\partial y} \kappa(\vec{r}), \quad \tau_{sc} = \frac{\partial^2}{\partial x \partial y} \chi$$

Local density approximation

$$\rho_{ss'}(\vec{r}_1, \vec{r}_2, t) = \rho_s(\vec{r}_1)\rho_{s'}(\vec{r}_2)(1 + d_{ss'}(\vec{r}_1, \vec{r}_2)) \quad s, s' \in \{+, -\}$$



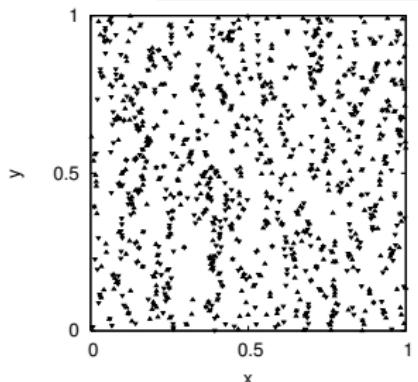
Relaxed dislocation configuration



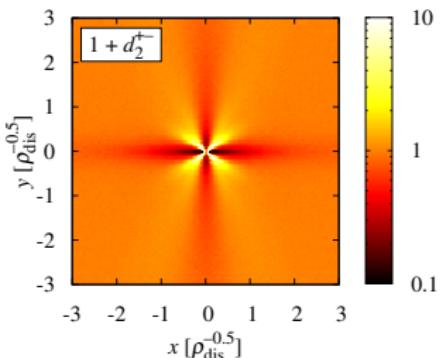
Correlation function

Local density approximation

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Relaxed dislocation configuration



Correlation function

$$d_{ss'}(\vec{r}_1, \vec{r}_2) = d_{ss'}(\vec{r}_1 - \vec{r}_2, \rho(\vec{r}_1))$$

Local density approximation

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Stress "like" terms with $\rho = \rho_+ + \rho_-$ and $\kappa = \rho_+ - \rho_-$

$$\begin{aligned}\partial_t \rho_+(\vec{r}, t) + M_0 b \partial_x \left\{ \rho_+ \left[\tau_{\text{ext}} + \tau_{\text{sc}} - \tau_f + \tau_b + \tau_a + \frac{\kappa}{\rho} \tau_f \right] \right\} &= 0 \\ \partial_t \rho_-(\vec{r}, t) - M_0 b \partial_x \left\{ \rho_- \left[\tau_{\text{ext}} + \tau_{\text{sc}} - \tau_f + \tau_b - \tau_a - \frac{\kappa}{\rho} \tau_f \right] \right\} &= 0\end{aligned}$$

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Correlation functions

$$\begin{aligned}d_s &= \frac{1}{2}(d_{++} + d_{--} + d_{+-} + d_{-+}), \\ d_p &= \frac{1}{2}(d_{++} + d_{--} - d_{+-} - d_{-+}), \\ d_a &= \frac{1}{2}(d_{++} - d_{--} - d_{+-} + d_{-+}), \\ d_{a'} &= \frac{1}{2}(d_{++} - d_{--} + d_{+-} - d_{-+}).\end{aligned}$$

Local density approximation

$$\begin{aligned}\rho_{ss'}(\vec{r}_1, \vec{r}_2, t) &= \rho_s(\vec{r}_1)\rho_{s'}(\vec{r}_2)(1 + d_{ss'}(\vec{r}_1, \vec{r}_2)) \quad s, s' \in \{+, -\} \\ d_{ss'}(\vec{r}_1, \vec{r}_2) &= d_{ss'}((\vec{r}_1 - \vec{r}_2)\sqrt{\rho(\vec{r}_1)})\end{aligned}$$

Stress “like” terms with $\rho = \rho_+ + \rho_-$ and $\kappa = \rho_+ - \rho_-$

$$\begin{aligned}\partial_t \rho_+(\vec{r}, t) + M_0 b \partial_x \left\{ \rho_+ \left[\tau_{\text{ext}} + \tau_{\text{sc}} - \tau_f + \tau_b + \tau_a + \frac{\kappa}{\rho} \tau_f \right] \right\} &= 0 \\ \partial_t \rho_-(\vec{r}, t) - M_0 b \partial_x \left\{ \rho_- \left[\tau_{\text{ext}} + \tau_{\text{sc}} - \tau_f + \tau_b - \tau_a - \frac{\kappa}{\rho} \tau_f \right] \right\} &= 0\end{aligned}$$

$$\begin{aligned}\tau_{\text{sc}}(\vec{r}) &= \int \tau_{\text{ind}}(\vec{r} - \vec{r}') \kappa(\vec{r}') d^2 r', \quad \frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_y \kappa \\ \tau_f(\vec{r}) &= - \int \rho(\vec{r}') d_a(\vec{r} - \vec{r}') \tau_{\text{ind}}(\vec{r} - \vec{r}') d^2 r' \\ \tau_b(\vec{r}) &= \int \kappa(\vec{r}') d_s(\vec{r} - \vec{r}') \tau_{\text{ind}}(\vec{r} - \vec{r}') d^2 r' \\ \tau_a(\vec{r}) &= \int \rho(\vec{r}') d_p(\vec{r} - \vec{r}') \tau_{\text{ind}}(\vec{r} - \vec{r}') d^2 r'\end{aligned}$$

Local density approximation

$$\begin{aligned}\rho_{ss'}(\vec{r}_1, \vec{r}_2, t) &= \rho_s(\vec{r}_1)\rho_{s'}(\vec{r}_2)(1 + d_{ss'}(\vec{r}_1, \vec{r}_2)) \quad s, s' \in \{+, -\} \\ d_{ss'}(\vec{r}_1, \vec{r}_2) &= d_{ss'}((\vec{r}_1 - \vec{r}_2)\sqrt{\rho(\vec{r}_1)})\end{aligned}$$

Stress “like” terms with $\rho = \rho_+ + \rho_-$ and $\kappa = \rho_+ - \rho_-$

$$\begin{aligned}\partial_t \rho_+(\vec{r}, t) + M_0 b \partial_x \left\{ \rho_+ \left[\tau_{\text{ext}} + \tau_{\text{sc}} - \tau_f + \tau_b + \tau_a + \frac{\kappa}{\rho} \tau_f \right] \right\} &= 0 \\ \partial_t \rho_-(\vec{r}, t) - M_0 b \partial_x \left\{ \rho_- \left[\tau_{\text{ext}} + \tau_{\text{sc}} - \tau_f + \tau_b - \tau_a - \frac{\kappa}{\rho} \tau_f \right] \right\} &= 0\end{aligned}$$

$$\begin{aligned}\tau_{\text{sc}}(\vec{r}) &= \int \tau_{\text{ind}}(\vec{r} - \vec{r}') \kappa(\vec{r}') d^2 r', \quad \frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_y \kappa \\ \tau_f(\vec{r}) &= \text{Friction stress } \mu b \alpha \sqrt{\rho(\vec{r})}, \\ \tau_b(\vec{r}) &= -Gb \frac{D}{\rho} \partial_x \kappa(\vec{r}), \\ \tau_a(\vec{r}) &= -Gb \frac{A}{\rho} \partial_x \rho(\vec{r}),\end{aligned}$$



Phase field theory

Balance equations for ρ_+ and ρ_-

$$\partial_t \rho_{\pm} + \partial_x [\rho_{\pm} v_{\pm}] = \pm f(\rho_+, \rho_-)$$



Phase field theory

Balance equations for ρ_+ and ρ_- with no sources

$$\partial_t \rho_{\pm} + \partial_x [\rho_{\pm} v_{\pm}] = 0$$



Phase field theory

Balance equations for ρ_+ and ρ_- with no sources

$$\partial_t \rho_{\pm} + \partial_x [\rho_{\pm} v_{\pm}] = 0$$

Dissipative motion ($v_{\pm} \propto \mp \tau$)

$$\tau = \frac{\delta E}{\delta \gamma}, \quad E \left[\kappa = \frac{1}{b} \partial_x \gamma \right] \Rightarrow \tau = -\partial_x \frac{\delta E}{\delta \kappa}, \Rightarrow v_{\pm} = -M_0 \partial_x \frac{\delta E}{\delta \kappa}$$

Phase field theory

Balance equations for ρ_+ and ρ_- with no sources

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Phase field functional $P[\rho_+, \rho_-, \chi]$

$$v_+ = -M_0 \left\{ \partial_x \left[\frac{1+\beta}{2} \frac{\delta P}{\delta \rho_+} - \frac{1-\beta}{2} \frac{\delta P}{\delta \rho_-} \right] \right\}$$

$$v_- = -M_0 \left\{ \partial_x \left[\frac{1+\beta}{2} \frac{\delta P}{\delta \rho_-} - \frac{1-\beta}{2} \frac{\delta P}{\delta \rho_+} \right] \right\}$$

Phase field theory

Balance equations for ρ_+ and ρ_- with no sources

$$\partial_t \rho_{\pm} + \partial_x [\rho_{\pm} v_{\pm}] = 0$$

Dissipative motion ($v_{\pm} \propto \mp \tau$)

$$\tau = \frac{\delta E}{\delta \gamma}, \quad E \left[\kappa = \frac{1}{b} \partial_x \gamma \right] \Rightarrow \tau = -\partial_x \frac{\delta E}{\delta \kappa}, \Rightarrow v_{\pm} = -M_0 \partial_x \frac{\delta E}{\delta \kappa}$$

Phase field functional $P[\rho_+, \rho_-, \chi]$

$$\begin{aligned} \partial_t \rho_+ - \partial_x \left\{ \rho_+ M_0 \left[\partial_x \frac{\delta P}{\delta \kappa} + \beta \partial_x \frac{\delta P}{\delta \rho} \right] \right\} &= 0 \\ \partial_t \rho_- + \partial_x \left\{ \rho_- M_0 \left[\partial_x \frac{\delta P}{\delta \kappa} - \beta \partial_x \frac{\delta P}{\delta \rho} \right] \right\} &= 0 \end{aligned}$$



Phase field functional

Coarse grained energy functional

$$P_{\text{sc}}[\chi, \rho_+, \rho_-] = \int \left[-\frac{1-\nu}{4\mu} (\Delta \chi)^2 + b\chi \partial_y \kappa \right] d^2 r$$



Phase field functional

Coarse grained energy functional

$$P_{\text{sc}}[\chi, \rho_+, \rho_-] = \int \left[-\frac{1-\nu}{4\mu} (\Delta \chi)^2 + b\chi \partial_y \kappa \right] d^2 r$$

Minimum condition

$$\frac{\delta P_{\text{sc}}}{\delta \chi} = 0$$

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_y \kappa, \quad \tau_{\text{sc}} = \partial_x \partial_y \chi$$

Phase field functional

Coarse grained energy functional

$$P_{\text{sc}}[\chi, \rho_+, \rho_-] = \int \left[-\frac{1-\nu}{4\mu} (\Delta \chi)^2 + b \chi \partial_y \kappa \right] d^2 r$$

Minimum condition

$$\begin{aligned} \frac{\delta P_{\text{sc}}}{\delta \chi} &= 0 \\ \frac{1-\nu}{2\mu} \Delta^2 \chi &= b \partial_y \kappa, \quad \tau_{\text{sc}} = \partial_x \partial_y \chi \end{aligned}$$

Evolution equation

$$\frac{\partial \rho(\vec{r}, t)}{\partial t} + M_0 \vec{b} \frac{\partial}{\partial \vec{r}} [\kappa(\vec{r}, t) \{\tau_{\text{sc}}(\vec{r}, t) + \tau_{\text{ext}}\}] = 0$$

$$\frac{\partial \kappa(\vec{r}, t)}{\partial t} + M_0 \vec{b} \frac{\partial}{\partial \vec{r}} [\rho(\vec{r}, t) \{\tau_{\text{sc}}(\vec{r}, t) + \tau_{\text{ext}}\}] = 0$$

Correlation part

Correlation contribution

$$P[\rho_+, \rho_-, \chi] = P_{\text{sc}}[\rho_+, \rho_-, \chi] + P_{\text{corr}}[\rho_+, \rho_-]$$

$$P_{\text{corr}}^{\pm}[\rho_{\pm}] = \int \left[GAb\rho \ln \left(\frac{\rho}{\rho_0} \right) + \frac{GDb}{2} \frac{\kappa^2}{\rho} \right] d^2r$$

Correlation part

Correlation contribution

$$P[\rho_+, \rho_-, \chi] = P_{\text{sc}}[\rho_+, \rho_-, \chi] + P_{\text{corr}}[\rho_+, \rho_-]$$

$$P_{\text{corr}}^{\pm}[\rho_{\pm}] = \int \left[GAb\rho \ln \left(\frac{\rho}{\rho_0} \right) + \frac{GDb}{2} \frac{\kappa^2}{\rho} \right] d^2r$$

Evolution equations

$$\partial_t \rho = -Mb\partial_x \left\{ \kappa\tau_{\text{sc}} - GbD\frac{\kappa}{\rho}\partial_x\kappa - GbA\partial_x\rho \right\}$$

$$\partial_t \kappa = -Mb\partial_x \left\{ \rho\tau_{\text{sc}} - GbD\partial_x\kappa - GbA\frac{\kappa}{\rho}\partial_x\rho \right\}$$



Flow stress

Adding flow stress

$$\partial_t \rho = -Mb\partial_x \left\{ \kappa \tau_{sc} - GbD \frac{\kappa}{\rho} \partial_x \kappa - GbA \partial_x \rho \right\}$$

$$\partial_t \kappa = -Mb\partial_x \left\{ \rho \tau_{sc} - \rho \left(1 - \frac{\kappa^2}{\rho^2} \right) \tau_f - GbD \partial_x \kappa - GbA \frac{\kappa}{\rho} \partial_x \rho \right\}$$

$$\tau_f(\tau_{sc}) = \begin{cases} \tau_{sc} & \text{if } \tau_{sc} < \alpha \mu b \sqrt{\rho} \\ \alpha \mu b \sqrt{\rho} & \text{if } \tau_{sc} > \alpha \mu b \sqrt{\rho} \end{cases}$$

Flow stress

Adding flow stress

$$\partial_t \rho = \partial_x \left\{ \kappa M_0 \partial_x \frac{\delta P}{\delta \kappa} + \rho M_0 \partial_x \frac{\delta P}{\delta \rho} \right\}$$

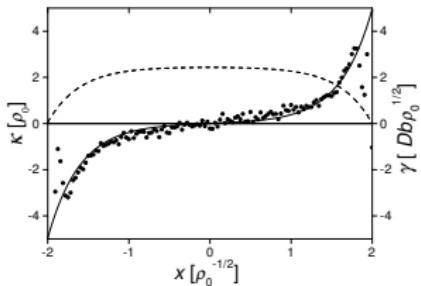
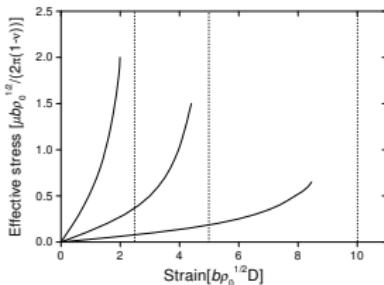
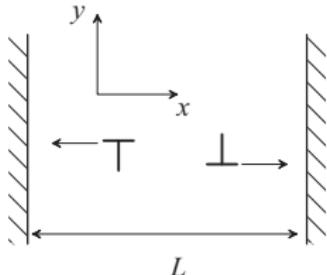
$$\partial_t \kappa = \partial_x \left\{ \rho \mathbf{M} \left(\partial_x \frac{\delta P}{\delta \kappa} \right) + \kappa M_0 \partial_x \frac{\delta P}{\delta \rho} \right\}$$

$$M(x) = M_0 \begin{cases} \frac{\kappa^2}{\rho^2} x & \text{if } |x| < x_0 \\ \operatorname{sgn}(x) \left[|x| - x_0 \left(1 - \frac{\kappa^2}{\rho^2} \right) \right] & \text{if } |x| \geq x_0 \end{cases}$$

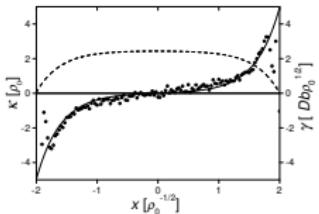
with $x_0 = \alpha \mu b^2 \sqrt{\rho}$



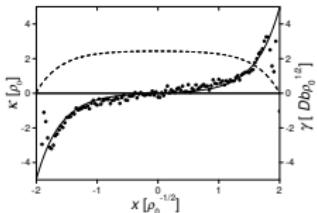
Finite size effect



Finite size effect



Finite size effect

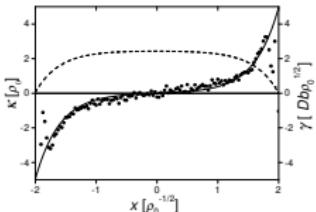


if $\rho = 1/l^2$

$$\tau = GbD \frac{1}{\rho(\vec{r})} \frac{\partial \kappa(\vec{r})}{\partial x}$$

$$\kappa(\vec{r}) = \tau \frac{1}{l^2} \frac{1}{GbD} x$$

Finite size effect

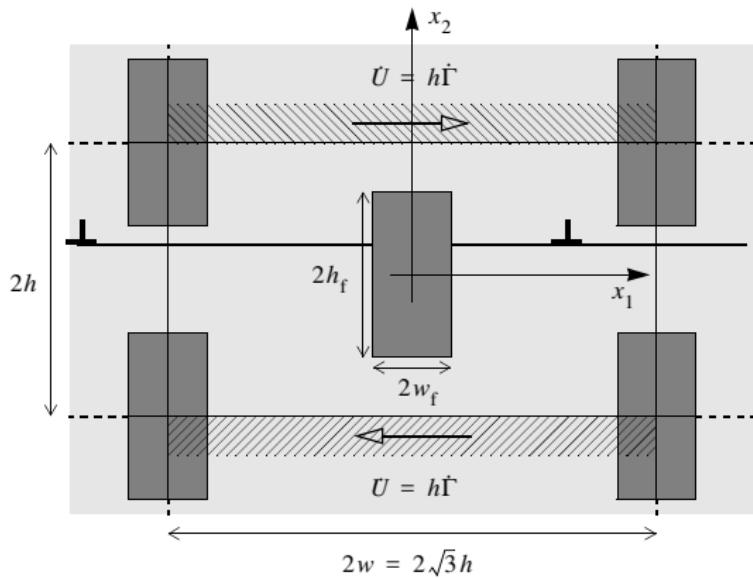


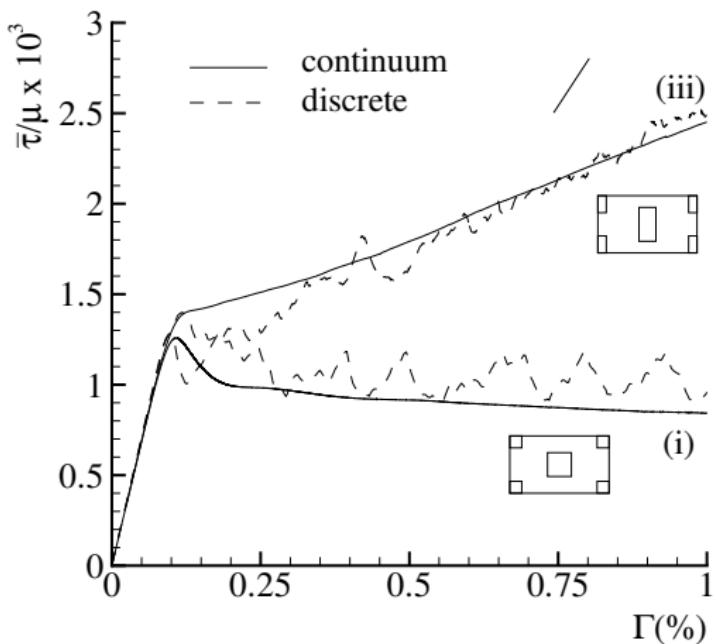
if $\rho = 1/l^2$

$$\begin{aligned}\tau &= GbD \frac{1}{\rho(\vec{r})} \frac{\partial \kappa(\vec{r})}{\partial x} \\ \kappa(\vec{r}) &= \tau \frac{1}{l^2} \frac{1}{GbD} x\end{aligned}$$

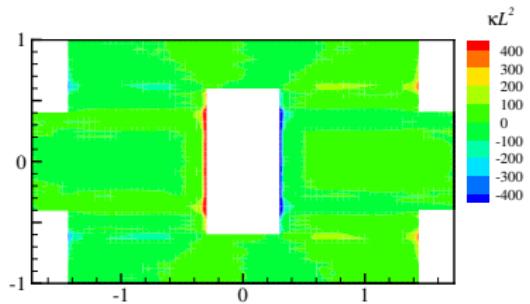
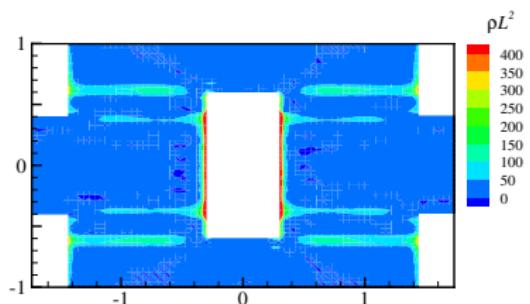
at the surface $-L/2 \leq \rho = -\kappa$

$$\begin{aligned}\tau &= -GbD \frac{1}{\kappa(\vec{r})} \frac{\partial \kappa(\vec{r})}{\partial x} \\ \kappa(\vec{r}) &= \kappa_0 \exp \left\{ -\tau \frac{1}{GbD} x \right\}\end{aligned}$$





NG composite





Debye screening

Single dislocation

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_2 \delta$$



Debye screening

Single dislocation+"bath"

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_2 \delta + \textcolor{blue}{b \partial_2 \kappa}$$



Debye screening

Single dislocation+"bath"

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_2 \delta + b \partial_2 \kappa$$

Single dislocation+induced GND

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_2 \left\{ \delta + \frac{\rho_0}{GbD} \partial_2 \chi \right\}$$



Debye screening

Single dislocation+"bath"

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_2 \delta + b \partial_2 \kappa$$

Single dislocation+induced GND

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b \partial_2 \left\{ \delta + \frac{\rho_0}{GbD} \partial_2 \chi \right\}$$

Solution

$$\chi \propto \sinh(k_0 y) K_0(k_0 r)$$

Debye screening

Single dislocation+"bath"

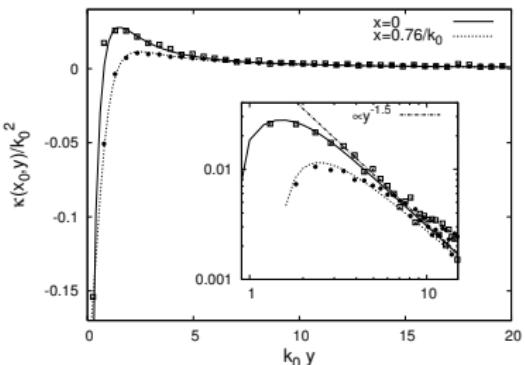
$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b\partial_2 \delta + b\partial_2 \kappa$$

Single dislocation+induced GND

$$\frac{1-\nu}{2\mu} \Delta^2 \chi = b\partial_2 \left\{ \delta + \frac{\rho_0}{GbD} \partial_2 \chi \right\}$$

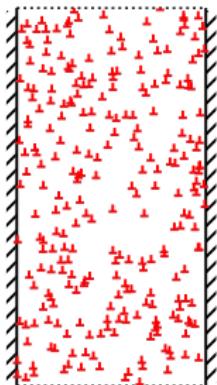
Solution

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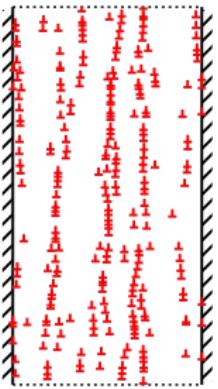


Only GND

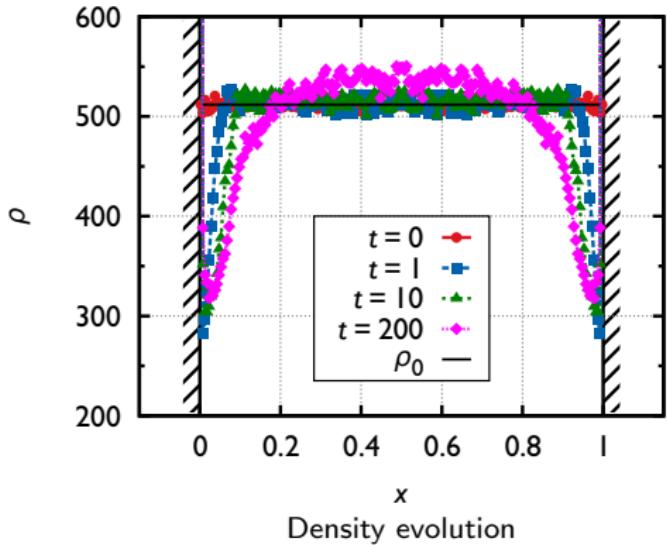




Only GND



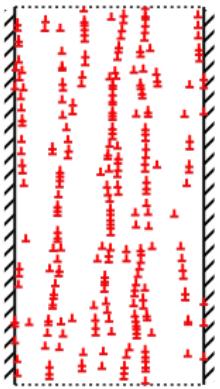
Dislocation configuration



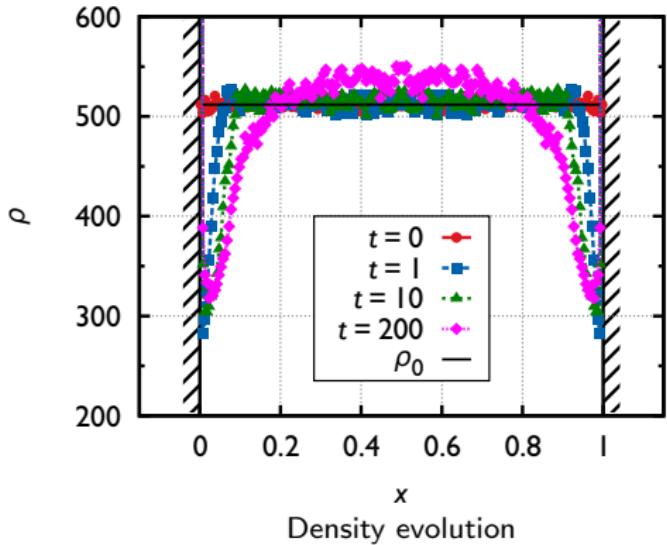
Density evolution



Only GND



Dislocation configuration



Density evolution

$$j_+ = -M_0 AbG \partial_x \rho_+ \equiv 0$$



Beyond Gradient term

Plastic potential

$$P = \int_G \left\{ -\frac{(1-\nu)}{4\mu} (\Delta \chi)^2 + b\chi \partial_y \rho_+ + GbA\rho_+ \ln(\rho_+/ \rho_0) + \rho_+ U [(\partial_x \rho_+)^2 / \rho_+^3] \right\} d^2r$$



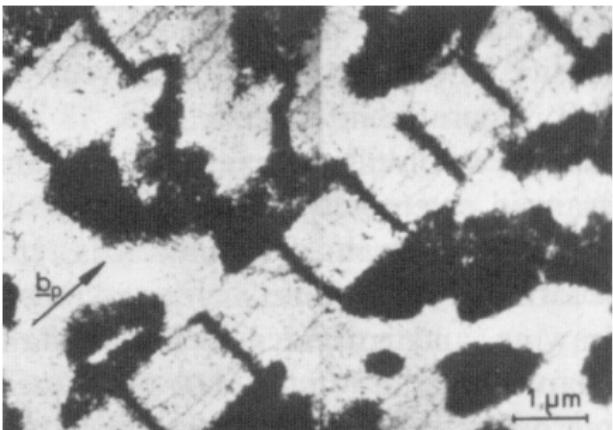
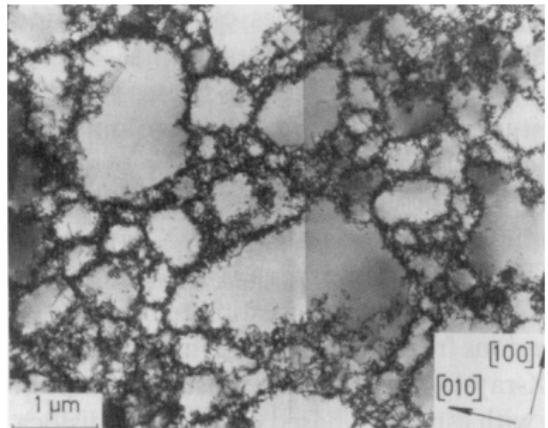
Beyond Gradient term

Surface term \Rightarrow Boundary condition

$$P_{surface} = \int_{\partial G} \lambda \sqrt{\rho_+} (\vec{b} \vec{n}) ds$$



Dislocation patterning



Stability analysis

Trivial solution

$$\begin{aligned}\rho(\vec{r}, t) &= \rho + \delta\rho(\vec{r}, t) \\ \kappa(\vec{r}, t) &= \delta\kappa(\vec{r}, t) \\ \chi(\vec{r}, t) &= \tau_0 xy + \delta\chi(\vec{r}, t)\end{aligned}$$

Periodic perturbation

$$\begin{pmatrix} \delta\rho \\ \delta\kappa \\ \delta\chi \end{pmatrix} = \begin{pmatrix} \delta\rho_0 \\ \delta\kappa_0 \\ \delta\chi_0 \end{pmatrix} \exp\left(\frac{\lambda}{t_0}t + i\sqrt{\rho_0}\vec{k}\vec{r}\right)$$

Stability analysis

Trivial solution

$$\begin{aligned}\rho(\vec{r}, t) &= \rho + \delta\rho(\vec{r}, t) \\ \kappa(\vec{r}, t) &= \delta\kappa(\vec{r}, t) \\ \chi(\vec{r}, t) &= \tau_0 xy + \delta\chi(\vec{r}, t)\end{aligned}$$

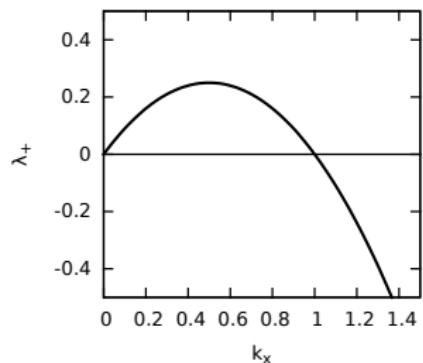
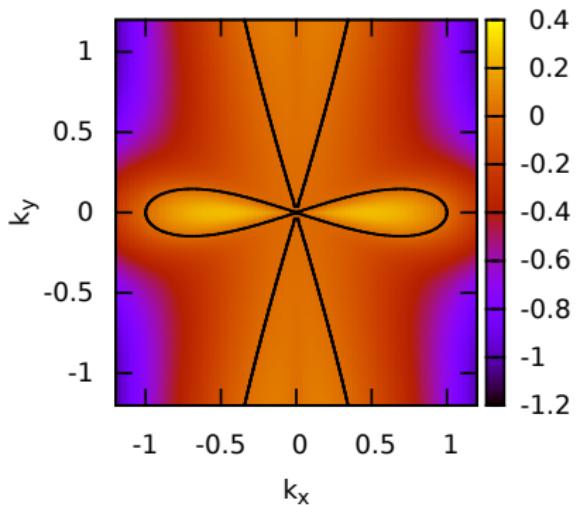
Periodic perturbation

$$\begin{pmatrix} \delta\rho \\ \delta\kappa \\ \delta\chi \end{pmatrix} = \begin{pmatrix} \delta\rho_0 \\ \delta\kappa_0 \\ \delta\chi_0 \end{pmatrix} \exp\left(\frac{\lambda}{t_0}t + i\sqrt{\rho_0}\vec{k}\vec{r}\right)$$

Stability condition

$$(\dot{\gamma}' + 2\alpha')(\dot{\gamma}' - \alpha') + A \frac{k_x^2 k_y^2}{k^4} + ADk_x^2 < 0.$$

Length scale selection





Michael Zaiser
Erik Van der Giessen
Peter Dusan Ispanovity
Geza Györgyi
Alphonse Finel

Serge Yefimov
Carmen Migual
Elias Aifantis
Thomas Hochrainer
Daniel Weygand
Stuart Pawley
Gilles Canova
Marc Verdier

Fields: $\rho(\vec{r}, \varphi)$ and $k(\vec{r}, \varphi)$

$$\hat{\nabla} = \rho(\vec{r}, \varphi) + \partial_\phi$$

$$L(\vec{r}, \varphi) = (\cos(\phi), \sin(\phi), 0, k(\vec{r}, \varphi))$$

Velocities

$$\vec{v} = (v_1, v_2, 0) = v(\vec{r}, \varphi)(\sin(\phi), -\cos(\phi), 0)$$

$$V(\vec{r}, \varphi) = (v_1, v_2, 0, -\nabla_L v)$$

Evolution eqs. $q = \rho k$

$$\partial_t \rho = -\hat{\nabla}(\rho V) + qv$$

$$\partial_t q = -\hat{\nabla}(qV) - \rho \hat{\nabla}_L \hat{\nabla}_L v$$